## ORGANOMETALLICS

# Observation of alkaline earth complexes $\mathrm{M}(\mathrm{CO})_{8}(\mathrm{M}=\mathrm{Ca}, \mathrm{Sr}$, or Ba$)$ that mimic transition metals 


#### Abstract

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The alkaline earth metals calcium (Ca), strontium (Sr), and barium (Ba) typically engage in chemical bonding as classical main-group elements through their ns and $n p$ valence orbitals, where $n$ is the principal quantum number. Here we report the isolation and spectroscopic characterization of eight-coordinate carbonyl complexes $M(C O)_{8}$ (where $\mathrm{M}=\mathrm{Ca}, \mathrm{Sr}$, or Ba ) in a low-temperature neon matrix. Analysis of the electronic structure of these cubic $O_{h}$-symmetric complexes reveals that the metal-carbon monoxide (CO) bonds arise mainly from $\left[\mathrm{M}\left(\mathrm{d}_{\pi}\right)\right] \rightarrow(\mathrm{CO})_{8} \pi$ backdonation, which explains the strong observed red shift of the $\mathrm{C}-\mathrm{O}$ stretching frequencies. The corresponding radical cation complexes were also prepared in gas phase and characterized by mass-selected infrared photodissociation spectroscopy, confirming adherence to the 18 -electron rule more conventionally associated with transition metal chemistry.


Ihe periodic table of the elements is conventionally divided according to the valence atomic orbitals (AOs) into main-group s and p blocks, a transition metal d block, and a lanthanide and actinide $f$ block. A useful set of guidelines for understanding the structures and stabilities of molecules encompasses the associated 8 -, 18 -, and 32 -electron rules introduced by Langmuir $(1,2)$ before the advent of quantum theory. These rules were later explained by attributing particular stability to filled sp , spd , or spdf valence shells, respectively (3).
The alkaline earth elements beryllium, magnesium, calcium, strontium, and barium have a $n \mathrm{~s}^{2}$ valence-shell configuration, where $n$ is the principal quantum number, and, as such, typically engage in chemical bonding as ionic salt compounds or in polar bonds via their two $n s$ valence electrons in divalent M(II) species (4), where M is an alkaline earth metal. Earlier studies suggested that the heaviest atom barium may use its 5 d AOs to some extent in chemical bonds (5), which led to the suggestion that barium be designated an "honorary transition metal" (6). Previously, we reported the experimental observation of barium carbonyl ions $\mathrm{Ba}(\mathrm{CO})^{q}$ (where charge $q=+1$ and -1 ) (7). The analysis of the electronic structure showed that the cation binds the ligand mainly through $\mathrm{Ba}^{+}\left(5 \mathrm{~d}_{\pi}{ }^{1}\right) \rightarrow$

[^0]CO( $\pi^{*}$ LUMO) backdonation (LUMO, lowest unoccupied molecular orbital), with $\mathrm{Ba}^{+}$in the excited ${ }^{2} \mathrm{D}\left(5 \mathrm{~d}^{1}\right)$ electronic reference state. In that respect, the $\mathrm{Ba}(\mathrm{CO})^{+}$complex behaves similarly to a transition metal carbonyl. The Ba-CO interactions in the radical anion $\mathrm{Ba}(\mathrm{CO})^{-}$were consistent with dominant contributions of Ba $\left(5 \mathrm{~d}_{\pi}\right) \leftarrow \mathrm{CO}^{-}\left(\pi^{*} \mathrm{SOMO}\right) \pi$ donation (SOMO, singly occupied molecular orbital) and $\mathrm{Ba}\left(5 \mathrm{~d}_{\sigma} / 6 \mathrm{~s}\right) \leftarrow$ $\mathrm{CO}^{-}(\sigma \mathrm{HOMO}) \sigma$ donation (HOMO, highest occupied molecular orbital). The most important valence functions of barium in $\mathrm{Ba}(\mathrm{CO})^{+}$cation and $\mathrm{Ba}(\mathrm{CO})^{-}$anion thus appeared to be the 5 d orbitals (7).

These findings inspired us to search for the 18 -electron octacarbonyl complex $\mathrm{Ba}(\mathrm{CO})_{8}$. Surprisingly, we found that not only barium but also the lighter homologs strontium and calcium formed octacarbonyl complexes $\mathrm{M}(\mathrm{CO})_{8}(\mathrm{M}=$ $\mathrm{Ca}, \mathrm{Sr}$, or Ba ) that can be stabilized in a lowtemperature neon matrix.

The neutral alkaline earth-carbonyl complexes were prepared by the reactions of pulsed laserevaporated metal atoms and carbon monoxide (CO) in solid neon and were investigated using Fourier transform infrared absorption spectroscopy. The experiments were carried out with a wide range of CO concentrations (from 0.02 to $2 \%$ relative to Ne on the basis of volume). In the experiments with relatively low CO concentrations, terminally bonded mononuclear low-coordinate carbonyl complexes with $\mathrm{C} \equiv \mathrm{O}$ stretching frequencies in the 2050 to $1800 \mathrm{~cm}^{-1}$ region were observed. Experiments with isotopically substituted CO samples allowed the unambiguous identification of some low-coordinate complexes through isotopic shifts and splittings. The barium di-, tri-, and tetracarbonyls can clearly be identified on the basis of spectra in the experiments with $0.03 \%{ }^{12} \mathrm{C}^{16} \mathrm{O}$ (fig. S1); $0.05 \%{ }^{12} \mathrm{CO}$ and $0.05 \%{ }^{13} \mathrm{CO}$ (fig. S2); and $0.05 \% \mathrm{C}^{16} \mathrm{O}$ and
$0.05 \% \mathrm{C}^{18} \mathrm{O}$ (fig. S3). The monocarbonyls were theoretically predicted to be unstable (8) and were not observed in the experimental vibrational spectra. Intense absorption bands centered at $1987 \mathrm{~cm}^{-1}$ for $\mathrm{Ca}, 1995 \mathrm{~cm}^{-1}$ for Sr , and $2014 \mathrm{~cm}^{-1}$ for Ba were observed upon progressive annealing of the samples to temperatures of 10 to 13 K under relatively high CO concentrations (Table 1). These absorptions become the dominant features in the spectra with high CO concentrations (see Fig. 1A for Ca and figs. S 4 and S 5 for Sr and Ba), suggesting that the absorber is the coordinatively saturated 18 -electron octacarbonyl complex. The observation of only one carbonyl stretching band suggests that these neutral octacarbonyls have the highest cubic $O_{h}$ symmetry. Experiments with mixtures of ${ }^{12} \mathrm{C}^{16} \mathrm{O}$ and ${ }^{13} \mathrm{C}^{16} \mathrm{O}$ and also ${ }^{12} \mathrm{C}^{16} \mathrm{O}$ and ${ }^{12} \mathrm{C}^{18} \mathrm{O}$ provided conclusive identification of these cubic octacarbonyl complexes. Although the bands of the Sr and Ba complexes are too broad to resolve isotopic splittings, the band of the Ca complex is sharp and intense in the spectra with relatively low CO concentrations (Fig. 1A); well-resolved mixed isotopic spectra could therefore be compared with calculations. The experimentally observed spectra are in good agreement with the simulated isotopic spectral features shown in figs. S 6 and S 7 .
The radical cations of the alkaline earthcarbonyl complexes were prepared in the gas phase by using a pulsed laser vaporization-supersonic-expansion ion source and studied by mass-selected infrared photodissociation spectroscopy in the carbonyl stretching-frequency region. Typical mass spectra are shown in figs. S8 to S10. Mononuclear metal-carbonyl cation complexes $\left[\mathrm{M}(\mathrm{CO})_{n}\right]^{+}(\mathrm{M}=\mathrm{Ca}, \mathrm{Sr}$, or Ba ) with $n$ as high as 10 to 15 were observed. These larger complexes contain both strongly bound CO ligands directly coordinated to the central metal ion and weakly bound, or tagged, CO ligands $(9,10)$. All of these complexes dissociated by elimination of a CO ligand after photoexcitation at the $\mathrm{C} \equiv \mathrm{O}$ stretching vibrational frequencies. The infrared spectra of $\left[\operatorname{Sr}(\mathrm{CO})_{n}\right]^{+}(n=6$ to 9$)$ are shown in Fig. 1B and the corresponding spectra of the Ca and Ba complexes in figs. S11 and S12. All the spectra for the $\left[\mathrm{M}(\mathrm{CO})_{n}\right]^{+}$complexes with $n=6$ to 8 feature a broad single band [full width at half maximum (FWHM) of $49 \mathrm{~cm}^{-1}$ for Ca , $52 \mathrm{~cm}^{-1}$ for Sr , and $29 \mathrm{~cm}^{-1}$ for Ba for the $n=8$ complexes] that is slightly red shifted relative to the free-CO absorption at $2143 \mathrm{~cm}^{-1}$. The dissociation efficiency increases substantially (about $70 \%$ for $\mathrm{Ca}, 220 \%$ for Sr , and $107 \%$ for Ba ) from $n=8$ to 9 , and the $n=9$ complexes have a much narrower bandwidth [FWHM of $36 \mathrm{~cm}^{-1}$ for $\mathrm{Ca}, 10 \mathrm{~cm}^{-1}$ for Sr , and $16 \mathrm{~cm}^{-1}$ for Ba]. Along with the intense band around $2100 \mathrm{~cm}^{-1}$, an additional weak band in the 2160 to $2180 \mathrm{~cm}^{-1}$ region is also observed for the $n=9$ complexes. The bands in this latter frequency region are assigned to the vibrations of weakly tagged external CO ligands ( 9,10 ). The appearance of the tagged CO band at $n=9$ indicates that the $n=8$ complexes are coordinatively saturated.
We carried out quantum chemical calculations using density functional theory and ab initio
methods to support the assignments of the vibrational spectra to the observed species and to examine the electronic structure of the carbonyl complexes. Figure 2A shows the optimized geometries of the neutral octacarbonyls calculated at the M06-2X-D3/def2-TZVPP level. The molecules have cubic $\left(O_{h}\right)$ symmetry and a triplet $\left({ }^{3} \mathrm{~A}_{1 \mathrm{~g}}\right)$ electronic ground state with the valence electron configuration $\mathrm{a}_{1 \mathrm{~g}}{ }^{2} \mathrm{t}_{1 \mathrm{u}}{ }^{6} \mathrm{t}_{2 \mathrm{~g}}{ }^{6} \mathrm{a}_{2 \mathrm{u}}{ }^{2} \mathrm{e}_{\mathrm{g}}{ }^{2}$. Calculations of the singlet state (fig. S13) gave structures with $D_{4 d}\left(\mathrm{Ca}\right.$ and Sr ) or $D_{4 h}$ symmetry (Ba), which were between 6.5 and $7.5 \mathrm{kcal} / \mathrm{mol}$ higher in energy than the corresponding triplet state species. Figure 2, B and C, shows the equilibrium geometries of the cations $\left[\mathrm{M}(\mathrm{CO})_{8}\right]^{+}$, which resemble the neutral species in the singlet state. Thus, $\left[\mathrm{Ca}(\mathrm{CO})_{8}\right]^{+}$and $\left[\mathrm{Sr}(\mathrm{CO})_{8}\right]^{+}$possess a $D_{4 d}$ structure and a ${ }^{2} \mathrm{~A}_{1}$ electronic ground state, whereas $\left[\mathrm{Ba}(\mathrm{CO})_{8}\right]^{+}$has $D_{4 h}$ symmetry and a ${ }^{2} \mathrm{~B}_{2 \mathrm{~g}}$. electronic state. Figure 2 also shows the calculated zero-point energy (ZPE)-corrected bond dissociation energies ( $D_{0}$ ) of the octacarbonyls for loss of one and eight CO ligand(s). At the M06-2X-D3/def2-TZVPP level, the $D_{0}$ values for the dissociation of one CO lie between $9.1 \mathrm{kcal} /$ mol for Ca and $11.5 \mathrm{kcal} / \mathrm{mol}$ for Sr for the $\mathrm{M}(\mathrm{CO})_{8}$ neutral complexes and between $8.4 \mathrm{kcal} / \mathrm{mol}$ for Ba and $9.6 \mathrm{kcal} / \mathrm{mol}$ for Sr for the $\left[\mathrm{M}(\mathrm{CO})_{8}\right]^{+}$ cation complexes. The theoretical $D_{0}$ values for loss of eight CO ligands yielding $\mathrm{M} / \mathrm{M}^{+}$in the electronic ground state were between $58.8 \mathrm{kcal} /$
mol for Sr and $63.3 \mathrm{kcal} / \mathrm{mol}$ for Ca for the neutrals and between $73.5 \mathrm{kcal} / \mathrm{mol}$ for Ba and $87.6 \mathrm{kcal} / \mathrm{mol}$ for Ca for the cations. The calculated values at the $\operatorname{CCSD}(\mathrm{T}) /$ def2-TZVPP level using the M06-2X-D3/def2-TZVPP optimized geometries were slightly smaller. The basis set superposition error is quite small ( 0.3 to $0.4 \mathrm{kcal} / \mathrm{mol}$ for single CO dissociation and 0.6 to $1.0 \mathrm{kcal} / \mathrm{mol}$ for loss of eight COs).

The theoretical wave numbers for the $\mathrm{C} \equiv \mathrm{O}$ stretching modes of $\mathrm{M}(\mathrm{CO})_{8}$ and $\left[\mathrm{M}(\mathrm{CO})_{8}\right]^{+}$are shown in Table 1 along with the experimental values. The calculated values refer to the harmonic antisymmetric stretching frequencies scaled by a factor of 0.941 , which comes from the ratio of the calculated stretching mode of free $\mathrm{CO}\left(2278 \mathrm{~cm}^{-1}\right)$ to the experimental value ( $2143 \mathrm{~cm}^{-1}$ ) (11). The calculated harmonic wave numbers are slightly higher than the experimental anharmonic values, but the trends for the different metals and for the isotopes of the neutral systems are in excellent agreement with the recorded values. The calculations suggest only one infrared (IR) active mode for the neutral complexes $\mathrm{M}(\mathrm{CO})_{8}$ and two closely lying IR active modes for the cations $\left[\mathrm{M}(\mathrm{CO})_{8}\right]^{+}$. The latter splitting is too small to be experimentally observed. Theory and experiment indicate a considerable red shift of the $\mathrm{C} \equiv \mathrm{O}$ frequency for the neutral systems and a much smaller red shift for the cations. The calculated red shifts and the isotopic variations of the cal-

culated and experimental wave numbers match each other.

Figure 3 shows the splitting of the valence orbitals of atoms with a spd valence shell in an octacoordinate cubic field with $O_{h}$ symmetry (12). The breakdown of the AOs into irreducible representations of the $O_{h}$ point group in an eight-coordinate ligand field is similar to the splitting in a six-coordinate octahedral field $(13,14)$, but there are two important differences. One concerns the splitting of the $(n-1) \mathrm{d}$ AOs of the metal. The degenerate $\mathrm{e}_{\mathrm{g}} \mathrm{AOs}$ in the octacoordinate cubic field are the $(n-1) \mathrm{d}_{\pi}$ AOs, and the triply degenerate $t_{2 g}$ AOs are the $(n-1) \mathrm{d}_{\sigma}$ AOs. By contrast, in the six-coordinate octahedral field, the degenerate $\mathrm{e}_{\mathrm{g}}$ AOs are the $(n-1) \mathrm{d}_{\sigma}$ AOs, and the triply degenerate $\mathrm{t}_{2 \mathrm{~g}}$ AOs have $(n-1) \mathrm{d}_{\pi}$ character. The second difference concerns the appearance of the $\mathrm{a}_{2 \mathrm{u}}$ molecular orbital (MO) in the cubic field (Fig. 3), which is absent in the octahedral field. The $\mathrm{a}_{2 \mathrm{u}} \mathrm{MO}$ is a ligand-only orbital, because there is no valence AO of the spd shell that possesses this symmetry. This has consequences for the number of electrons that are required to fulfill the 18 -electron rule. Because two valence electrons of the ligands in a cubic field are not available for donation to the central metal atom M in $\mathrm{ML}_{8}$, where L is a ligand, the complex must provide a total of 20 electrons to fully occupy the metal's valence shell. This scenario has recently been explored in the transition

Fig. 1. IR spectra of alkaline earth carbonyl
complexes. (A) IR absorption spectra of calcium-carbonyl complexes in the 2150 to $1850 \mathrm{~cm}^{-1}$ region from codeposition of laser-evaporated calcium atoms with $0.1 \% \mathrm{CO}$ in neon. Spectral lines: (i) after 30 min of sample deposition at 4 K , (ii) after a $12-\mathrm{K}$ annealing period, (iii) after a $13-\mathrm{K}$ annealing period, (iv) after 15 min of visible light irradiation, and (v) after another $12-\mathrm{K}$ annealing period. (B) IR photodissociation spectra of the $\operatorname{Sr}(\mathrm{CO})_{n}{ }^{+}(n=6$ to 9$)$ complexes in the 2300 to $1800 \mathrm{~cm}^{-1}$ region.

Fig. 2. Calculated equilibrium geometries of alkaline earth octacarbonyls. (A) $\mathrm{M}(\mathrm{CO})_{8}$ (M = Ca, Sr, or Ba), (B) $\left[\mathrm{M}(\mathrm{CO})_{8}\right]^{+}$ $(\mathrm{M}=\mathrm{Ca}$ or Sr$)$, and $(\mathbf{C})\left[\mathrm{Ba}(\mathrm{CO})_{8}\right]^{+}$ Bond lengths are in angstroms. The $D_{0}$ values in roman type are the ZPE-corrected bond dissociation energies for loss of one CO ligand; the italicized values are the corresponding energies for the loss of eight CO ligands and $\mathrm{M} / \mathrm{M}^{+}$in the ground state. The values without parentheses are from M06-2X-D3/def2-TZVPP calculations; the values in parentheses are from $\operatorname{CCSD}(\mathrm{T}) /$ def2-TZVPP using M06-2X-D3/def2-TZVPP optimized geometries.
metal octacarbonyl anions $\left[\mathrm{TM}(\mathrm{CO})_{8}\right]^{-}(\mathrm{TM}=$ $\mathrm{Sc}, \mathrm{Y}$, or La) (15). Only 18 electrons are available in $\mathrm{M}(\mathrm{CO})_{8}(\mathrm{M}=\mathrm{Ca}, \mathrm{Sr}$, or Ba$)$, so the degenerate $\mathrm{e}_{\mathrm{g}} \mathrm{MO}$ is occupied by two electrons with the same spin, giving a triplet $\left({ }^{3} \mathrm{~A}_{1 g}\right)$ electronic ground state. Because the $\mathrm{e}_{\mathrm{g}}$ correlates with the $(n-1) \mathrm{d}_{\pi}$ AO, it becomes clear that the electronic reference state of the alkaline earth atom in $\mathrm{M}(\mathrm{CO})_{8}$ is a triplet state with $n s^{0}(n-1) \mathrm{d}^{2}$ electron configuration. The metal center has a zero formal oxidation state.
We analyzed the nature of the metal-CO bonds with the energy decomposition analysis-natural orbitals for chemical valence (EDA-NOCV) (16) method, a powerful tool that provides detailed insight into chemical bonding (17). A description is given in the methods section. Table 2 shows the calculated results for the interactions between the metal atom M in the triplet electronic
reference state with a $n s^{0}(n-1) \mathrm{d}^{2}$ electron configuration and the $(\mathrm{CO})_{8}$ cage at the frozen geometry of $\mathrm{M}(\mathrm{CO})_{8}$. The interaction energies $\Delta E_{\text {int }}$ suggest that the intrinsic attraction $\mathrm{M}-(\mathrm{CO})_{8}$ is strong and varies in the order $\mathrm{Ca}>\mathrm{Sr} \gg \mathrm{Ba}$. The dominant contribution to the total interaction $\Delta E_{\text {int }}$ comes from the orbital term $\Delta E_{\text {orb. }}$. The preparation energies ( $\Delta E_{\text {prep }}$ ) involved in the formation of the $(\mathrm{CO})_{8}$ cage from free CO molecules are low, ranging from $3.3 \mathrm{kcal} / \mathrm{mol}$ for Ba to $13.9 \mathrm{kcal} /$ mol for Ca , whereas the electronic excitation energies for the atoms into the spherically symmetric $n \mathrm{~s}^{2} \rightarrow n \mathrm{~s}^{0}(n-1) \mathrm{d}^{2}$ triplet state of M are quite high, lying between $68.2 \mathrm{kcal} / \mathrm{mol}$ for Ba and $159.5 \mathrm{kcal} / \mathrm{mol}$ for Ca. The strongly stabilizing interaction energies $\Delta E_{\text {int }}$ caused by the attraction between the CO ligands and electronically excited M overcompensate the $\Delta E_{\text {prep }}$ values. The adiabatic interaction energy $\Delta E\left(=-D_{\mathrm{e}}\right.$,


Fig. 3. Bonding scheme and shape of the occupied valence orbitals of $\mathbf{M}(\mathrm{CO})_{8}(\mathbf{M}=\mathbf{C a}, \mathrm{Sr}$, or Ba). Splitting of the spd valence orbitals of an atom $M$ with the configuration $(n-1) \mathrm{d}^{2} n s^{0} n p^{0}$ in the octacoordinate cubic $\left(O_{h}\right)$ field of eight CO ligands is also given. Only the occupied valence orbitals that are relevant for the M-CO interactions are shown. Up and down arrows indicate electrons with opposite spin.
where $D_{\mathrm{e}}$ is the dissociation energy without zero-point vibrational energy correction) with respect to the electronic ground state of M and eight CO is between $-65.5 \mathrm{kcal} / \mathrm{mol}$ for Sr and $-73.7 \mathrm{kcal} / \mathrm{mol}$ for Ba. The $D_{\mathrm{e}}$ values in Table 2 exhibit a similar trend as the $D_{0}$ data in Fig. 2. The former values do not consider ZPE contributions, and they are obtained from calculations with different basis sets using Slater-type orbital basis functions.
The most important insight from the EDANOCV calculations is the breakdown of $\Delta E_{\text {orb }}$ into pairwise orbital interactions. Table 2 shows that the metal-CO bonding comes mainly from the $\left[\mathrm{M}\left(\mathrm{d}_{\pi}\right)\right] \rightarrow(\mathrm{CO})_{8} \pi$ backdonation of the degenerate ( $\mathrm{e}_{\mathrm{g}}$ ) set of singly occupied ( $n-1$ )d AOs of the metal into the antibonding $\pi^{*}$ MOs of CO . This explains the large red shift of the $\mathrm{C}=\mathrm{O}$ stretching mode of the octacarbonyls. The contribution of the $[\mathrm{M}(\mathrm{x})] \leftarrow(\mathrm{CO})_{8} \sigma$ donation into x , where x denotes the valence acceptor AO of M , is much smaller than the $\left[\mathrm{M}\left(\mathrm{d}_{\pi}\right)\right] \rightarrow(\mathrm{CO})_{8}$ $\pi$ backdonation. The order of the acceptor AOs of the metal atoms for the interaction energy is $(n-1) \mathrm{d}_{\sigma}>n \mathrm{~S}>n \mathrm{p}$. There is also a small stabilizing contribution of the $\mathrm{a}_{2 \mathrm{u}} \mathrm{MO}$, which is due to the polarization of the $(\mathrm{CO})_{8}$ ligand orbitals.

The dominating orbital interactions by the valence d electrons of $M$ can be explained with the energetically very-high-lying occupied orbitals, which make M atoms excellent donor species. In our previous study of $\mathrm{Ba}(\mathrm{CO})^{+}$, we found that the cation $\mathrm{Ba}^{+}\left(5 \mathrm{~d}_{\pi}^{1}\right)$ is a donor to neutral CO; the atomic partial charge of Ba in $\mathrm{Ba}(\mathrm{CO})^{+}$was calculated as $+1.39 e$. The experimental values for the energetically lowest-lying ${ }^{3} \mathrm{~F}$ state with the electron configuration $n s^{0}(n-1) \mathrm{d}^{2}$ are not far from the ionization limit (18). The valence d electrons of the M atoms are only weakly bonded to the atoms.

Figure 3 also shows the occupied valence MOs of $\mathrm{Ca}(\mathrm{CO})_{8}$ that are relevant for the metal-CO bonding. The shape of the five (19) MOs reveals that the contributions of the metal AOs are very small, except in the SOMO, which clearly exhibits the shape of the $(n-1) \mathrm{d}_{\pi}$ AOs. The valence MOs of

Table 1. Calculated and experimental wave numbers. Calculated (MO6-2X-D3/def2-TZVPP) and experimental wave numbers ( $\mathrm{cm}^{-1}$ ) of the $\mathrm{C}-0$ stretching mode and frequency shifts of $\mathrm{M}(\mathrm{CO})_{8}$ and $\left[\mathrm{M}(\mathrm{CO})_{8}\right]^{+}(\mathrm{M}=\mathrm{Ca}, \mathrm{Sr}$, or Ba$)$. The calculated values are scaled by 0.941 . Blank spaces indicate that isotope values are not available.

| Complex | Calculated wave numbers |  |  |  |  |  | Experimental wave numbers |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | co | $\Delta^{*}$ | ${ }^{13} \mathrm{CO}$ | $\Delta \dagger$ | $\mathrm{C}^{18} \mathrm{O}$ | $\Delta \dagger$ | co | $\Delta^{*}$ | ${ }^{13} \mathrm{CO}$ | $\Delta \dagger$ | $\mathrm{C}^{18} \mathrm{O}$ | $\Delta \dagger$ |
| $\mathrm{Ca}(\mathrm{CO})_{8}$ | 2018 | -125 | 1974 | -44 | 1969 | -49 | 1987 | -156 | 1944 | -43 | 1939 | -48 |
| $\mathrm{Sr}(\mathrm{CO})_{8}$ | 2027 | -116 | 1982 | -45 | 1977 | -50 | 1995 | -148 | 1952 | -43 | 1946 | -49 |
| $\mathrm{Ba}(\mathrm{CO})_{8}$ | 2044 | -99 | 1999 | -45 | 1994 | $-50$ | 2014 | -129 | 1970 | -44 | 1965 | -49 |
| $\left[\mathrm{Ca}(\mathrm{CO})_{8}\right]^{+}$ | $\begin{aligned} & 2116 \\ & 2117 \end{aligned}$ | $\begin{aligned} & -27 \\ & -26 \end{aligned}$ |  |  |  |  | 2094 | -49 |  |  |  |  |
| $\left[\mathrm{Sr}(\mathrm{CO})_{8}\right]^{+}$ | $\begin{aligned} & 2113 \\ & 2119 \end{aligned}$ | -30 -24 |  |  |  |  | 2096 | -47 |  |  |  |  |
| $\left[\mathrm{Ba}(\mathrm{CO})_{8}\right]^{+}$ | $\begin{aligned} & 2127 \\ & 2129 \end{aligned}$ | $\begin{aligned} & -16 \\ & -14 \end{aligned}$ |  |  |  |  | 2113 | -30 |  |  |  |  |

Table 2. EDA-NOCV results. EDA-NOCV results for triplet state $\mathrm{M}(\mathrm{CO})_{8}(\mathrm{M}=\mathrm{Ca}$, Sr , or Ba ) complexes at the M06-2X/TZ2P-ZORA level using M06-2X-D3/ def2-TZVPP optimized geometries, taking (CO) 8 in singlet ground state and M in triplet excited state with a $n s^{0}(n-1) d^{2}$ valence electronic configuration as interacting fragments. Energy values are given in kcal/mol.

| Energy term | Assignment | Interacting fragments |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | $\mathrm{Ca}+(\mathrm{CO})_{8}$ | $\mathrm{Sr}+(\mathrm{CO})_{8}$ | $\mathrm{Ba}+(\mathrm{CO})_{8}$ |
| $\Delta E_{\text {int }}$ |  | -243.1 | -224.1 | -145.2 |
| $\Delta E_{\text {hybria }}{ }^{*}$ |  | 41.8 | 46.9 | 37.4 |
| $\Delta E_{\text {Pauli }}$ |  | 19.5 | 30.5 | 30.8 |
| $\Delta E_{\text {elstat }} \dagger$ |  | -65.3 (21.5\%) | -61.7 (20.5\%) | -78.5 (36.8\%) |
| $\Delta E_{\text {orb }} \dagger$ |  | -239.1 (78.5\%) | -239.7 (79.5\%) | -135.0 (63.2\%) |
| $\Delta E_{\text {orb(1) }} \ddagger \S\left(\mathrm{e}_{\mathrm{g}}\right)$ | $[\mathrm{M}(\mathrm{d})] \rightarrow(\mathrm{CO})_{8} \pi$ backdonation | -206.2 (86.2\%) | -206.4 (86.2\%) | -95.0 (69.8\%) |
| $\Delta E_{\text {orb }}(2) \ddagger \S\left(\mathrm{t}_{2 \mathrm{~g}}\right)$ | $[\mathrm{M}(\mathrm{d})] \leftarrow(\mathrm{CO})_{8} \sigma$ donation | -21.3 (9.0\%) | -20.7 (8.7\%) | -22.8 (16.8\%) |
| $\Delta E_{\text {orb }}(3) \ddagger\left(\mathrm{a}_{1 \mathrm{~g}}\right)$ | $[\mathrm{M}(\mathrm{s})] \leftarrow(\mathrm{CO})_{8} \sigma$ donation | -2.4 (1.0\%) | -2.9 (1.2\%) | -3.2 (2.4\%) |
| $\Delta E_{\text {orb(4) }} \ddagger \oint^{( }\left(\mathrm{t}_{1 \mathrm{u}}\right)$ | $[\mathrm{M}(\mathrm{p})] \leftarrow(\mathrm{CO})_{8} \sigma$ donation | -0.9 (0.3\%) | -0.9 (0.3\%) | -2.7 (2.1\%) |
| $\Delta E_{\text {orb }}(5) \ddagger\left(\mathrm{a}_{2 \mathrm{u}}\right)$ | $(\mathrm{CO})_{8}$ polarization | -0.6 (0.3\%) | -1.0 (0.4\%) | -2.3 (1.7\%) |
| $\Delta E_{\text {orb }}$ (rest) $\ddagger$ |  | -7.7 (3.2\%) | -7.8(3.3\%) | -9.0 (6.7\%) |
| $\Delta E_{\text {prep(a) }}$ | $8 \mathrm{CO} \rightarrow(\mathrm{CO})_{8}$ | 13.9 | 7.7 | 3.3 |
| $\Delta E_{\text {prep(b) }} \\|$ | $\mathrm{M}, n \mathrm{~s}^{2} \rightarrow n s^{0}(n-1) \mathrm{d}^{2}(\mathrm{~T})$ | 159.5 | 150.9 | 68.2 |
| $\Delta E\left(\Delta E_{\text {int }}+\Delta E_{\text {prep }}\right)=-D_{\text {e }}$ | $\mathrm{M}(\mathrm{CO})_{8} \rightarrow \mathrm{M}\left({ }^{1} \mathrm{~S}\right)+8 \mathrm{CO}$ | -69.7 | -65.5 | -73.7 |

*Contribution of the metahybrid term in M06-2X. †The values in parentheses show the contribution to the total attractive interactions $\Delta E_{\text {elstat }}$ plus $\Delta E_{\text {orb }}$, where
 or three ( $t_{2 g}$ or $t_{1 u}$ ) components is given. \|The EDA calculations give a triplet state with spherically symmetrical distribution of the d electrons. The experimental values for excitation into the energetically lowest-lying ${ }^{3} \mathrm{~F}$ state with $n \mathrm{~s}^{\circ}(n-1) \mathrm{d}^{2}$ configuration are 124.2 and $59.8 \mathrm{kcal} / \mathrm{mol}$ for Ca and Ba , respectively. There is no experimental value for the relevant ${ }^{3} \mathrm{~F}$ state of Sr . The data are taken from (18).
the heavier homologs $\mathrm{Sr}(\mathrm{CO})_{8}$ and $\mathrm{Ba}(\mathrm{CO})_{8}$ look very similar (see figs. S14 and S15). The effect of the orbital interactions on the charge distribution is evident from the shape of the deformation densities $\Delta \rho$, which are associated with the orbital interactions. Figure S16 shows the contour line plots of the deformation densities $\Delta \rho_{(1)}$ to $\Delta \rho_{(5)}$, which are connected to the pairwise interactions $\Delta E_{\text {orb(1) }}$ to $\Delta E_{\text {orb(5) }}$ in $\mathrm{Ca}(\mathrm{CO})_{8}$ (Table 2). Note that only one component of the orbital terms is shown, and the color code of the charge flow is from red to blue. Figure S16A displays a large charge flow in the direction $\left[\mathrm{Ca}\left(\mathrm{d}_{\pi}\right)\right] \rightarrow$ $(\mathrm{CO})_{8}$, which comes from the $\pi$ backdonation. Figure S16, B to D, exhibits the charge flow in the opposite direction $[\mathrm{Ca}(\mathrm{x})] \leftarrow(\mathrm{CO})_{8}$ into the valence AOs of calcium $(n-1) \mathrm{d}_{\sigma}, n \mathrm{~s}$, and $n \mathrm{p}$. Figure S16E shows the charge polarization of the (CO) ${ }_{8} \mathrm{a}_{2 \mathrm{u}}$ ligand orbital, where electronic charge is shifted from oxygen toward carbon. The deformation densities $\Delta \rho_{(1)}$ to $\Delta \rho_{(5)}$ of the orbital interactions in $\mathrm{Sr}(\mathrm{CO})_{8}$ and $\mathrm{Ba}(\mathrm{CO})_{8}$ are shown in figs. S17 and S18, respectively.
The analysis of the electronic structure of the octacarbonyls $\mathrm{M}(\mathrm{CO})_{8}$ provides a straightforward explanation for why the molecules have a cubic $\left(O_{h}\right)$ equilibrium geometry and a triplet $\left({ }^{3} \mathrm{~A}_{1 g}\right)$ electronic ground state with a triplet reference state and $n s^{0}(n-1) \mathrm{d}^{2}$ electron configuration. The coordination by eight CO ligands in a cubic field fills the spd valence AOs of the alkaline earth atoms and the $\mathrm{a}_{2 \mathrm{u}}$ ligand-only MO. The energetically highest-lying bonding MO is the degenerate $\mathrm{e}_{\mathrm{g}}$ orbital. Because 16 electrons from the ligands and 2 electrons from metal are available, the $\mathrm{e}_{\mathrm{g}}$ MO has SOMO com-
ponents following Hund's rule. Simple electron counting indicates that the octacarbonyls $\mathrm{M}(\mathrm{CO})_{8}$ fulfill the 18 -electron rule. The EDA-NOCV calculations of the cations $\left[\mathrm{M}(\mathrm{CO})_{8}\right]^{+}$agree with the bonding analysis of the neutral complexes. The $\left[\mathrm{M}\left(\mathrm{d}_{\pi}\right)\right]^{+} \rightarrow(\mathrm{CO})_{8} \pi$ backdonation is weaker than in the neutral species, because there is only one $(n-1) \mathrm{d}_{\pi}$ electron available (tables S1 and S2).

The bonding situation in the alkaline earthoctacarbonyl complexes shows that not only barium but also strontium and calcium may effectively use their $(n-1) \mathrm{d}$ AOs in chemical bonding. It is conceivable that the chemical reactivity of heavier alkaline earth elements is more diverse than hitherto thought. Recent reports about unusual structures and reactivities of calcium and strontium compounds $(20,21)$ could be a hint that d-orbital participation of the alkaline earth metals plays an important role. The design of future experiments should consider the capacity of the heavier alkaline earth elements to behave like transition metals.

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## SUPPLEMENTARY MATERIALS

www.sciencemag.org/content/361/6405/912/suppl/DC1 Materials and Methods

Figs. S1 to S19
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Data S1
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## Science

# Observation of alkaline earth complexes $\mathrm{M}(\mathrm{CO})_{8}(\mathrm{M}=\mathrm{Ca}, \mathrm{Sr}$, or Ba$)$ that mimic transition metals 

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## Carbonyls in the s block

Conventional wisdom in chemistry distinguishes transition metals from other elements by their use of d orbitals in bonding. Wu et al. now report that alkaline earth metals can slide their electrons from s - to d -orbital bonding motifs as well (see the Perspective by Armentrout). Calcium, strontium, and barium all form coordination complexes with a cubic arrangement of eight carbonyl ligands and an 18 -electron valence shell. The compounds were characterized in frozen neon matrices by vibrational spectroscopy and in gas phase by mass spectrometry.

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